

# Finite element analysis for electronics

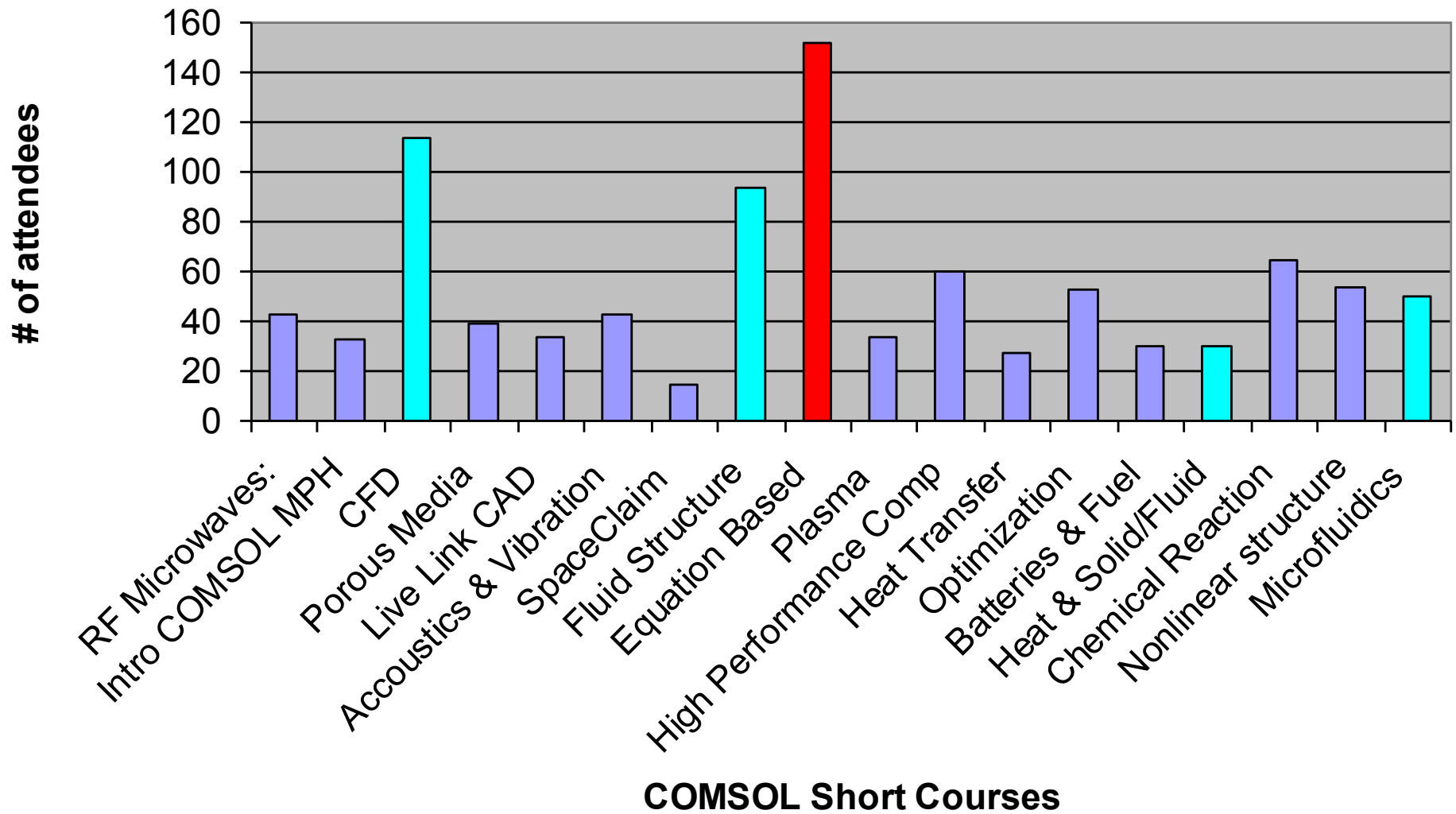
*October 8, 2010*

*COMSOL Conference 2010 Boston*

Yosuke Mizuyama, Ph.D.  
Panasonic Boston Laboratory

**Panasonic ideas for life**

## # of COMSOL short courses attendees



# Outline

- *Inkjet printhead*  
*Piezoelectric + free surface fluid dynamics coupling*
- *Grayscale photolithography for optical pickup*  
*Irreversible thermal process computation using COMSOL+MATLAB*
- *Holographic data storage*  
*Electromagnetic for infinite region*

# Electronics from Panasonic

## Home appliances

VT2 シリーズ 3Dビエラ誕生!

新しい世界を見よう!

3D VIERA FULL BLACK PANEL

見るテレビから「体験するテレビ」へ。迫力の3D対応ビエラ最高画質モデル

NEW

地上・BS・110度CSデジタルハイビジョンプラステレビ

65v型	TH-P65VT2
58v型	TH-P58VT2
54v型	TH-P54VT2
50v型	TH-P50VT2
46v型	TH-P46VT2
42v型	TH-P42VT2



<http://panasonic.jp>

## Laptop & cell phone

防水Wオープンスタイル®ゲートウェイモバイル

作成に役立つモバイル

14.1型 WXGA+

NEW

Windows 7 Professional

Intel CORE i3

変化した管理機能

ビジネスを快適に

Wi-Fi AX CERTIFIED

## Automotive



## Lighting & Health



## Energy

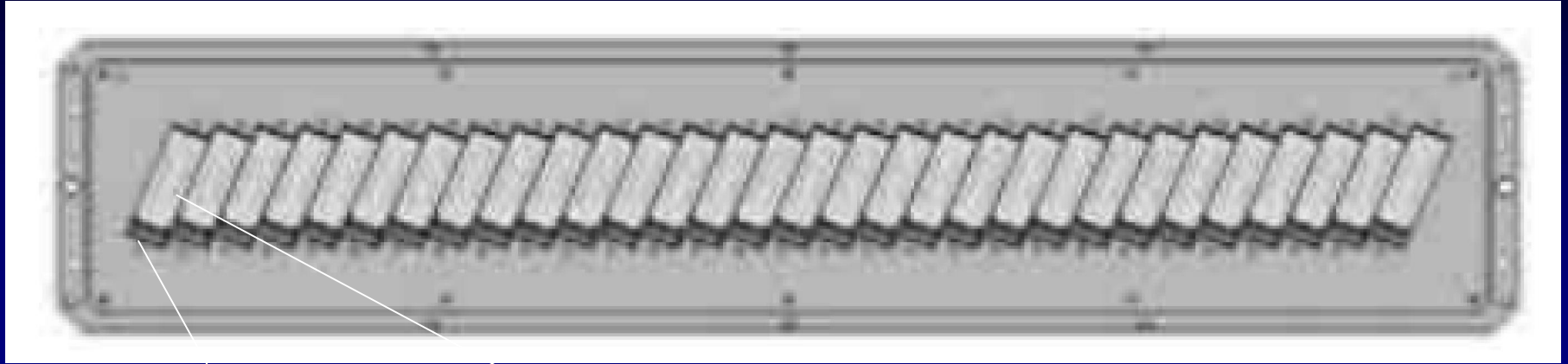


# Industrial inkjet printhead

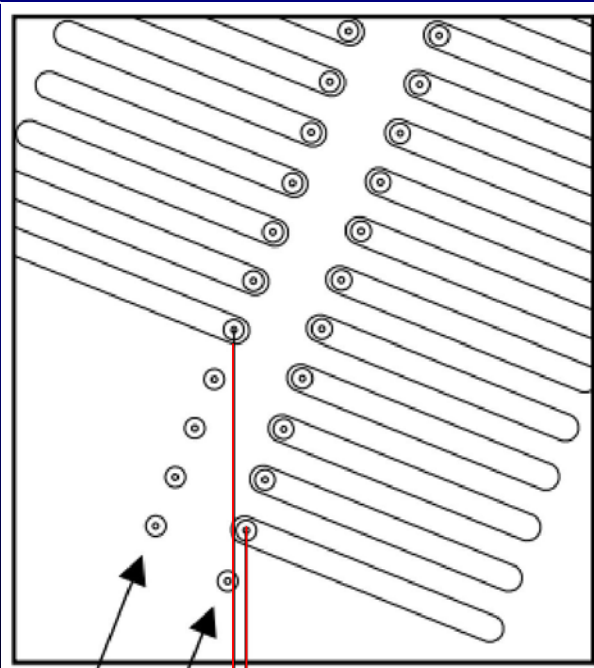


- 48,000 Jets
- 600 dpi
- 19.5" Width
- CMYK
- 240 ppm
- Aqueous Inks

# Industrial inkjet printhead



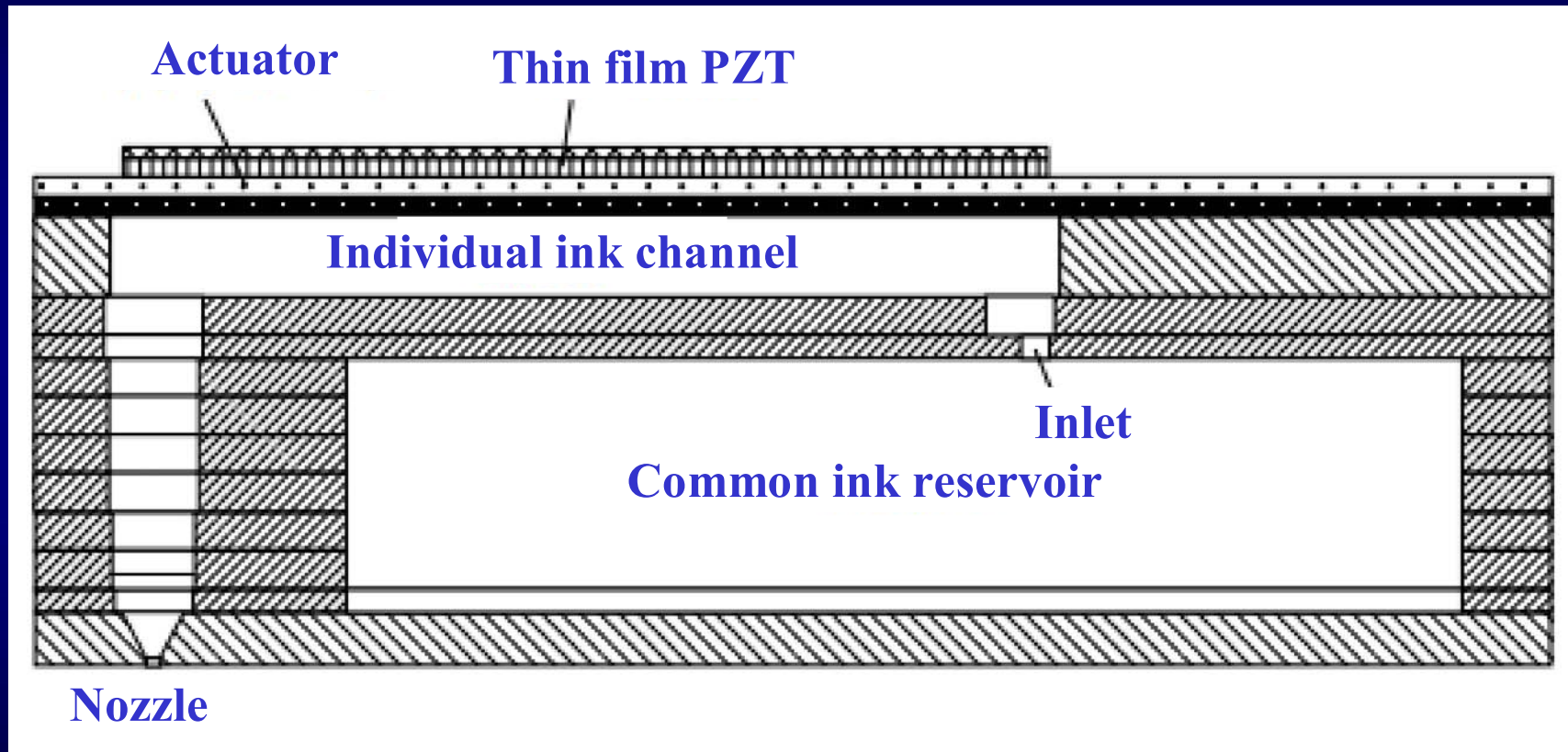
Paper feeding  
direction



600 dpi = 42.3um



# Industrial inkjet printhead



Cross section of the head

# Principle of PZT inkjet

Compliance Piezoelectric coupling

Strain  
Charge density  
displacement

$$\begin{matrix} 6 \times 1 \\ 3 \times 1 \end{matrix} \begin{bmatrix} S \\ D \end{bmatrix} = \begin{matrix} 6 \times 6 & 6 \times 3 \\ d & \epsilon_T \end{matrix} \begin{bmatrix} T \\ E \end{bmatrix} \begin{matrix} 6 \times 1 \\ 3 \times 1 \end{matrix}$$

Stress  
Electric field

Permittivity

d-matrix

$$d = \begin{bmatrix} 0 & 0 & 0 & 0 & * & 0 \\ 0 & 0 & 0 & * & 0 & 0 \\ d_{31} & d_{32} & d_{33} & 0 & 0 & 0 \end{bmatrix}$$

Strain-charge form  
constitutive equation



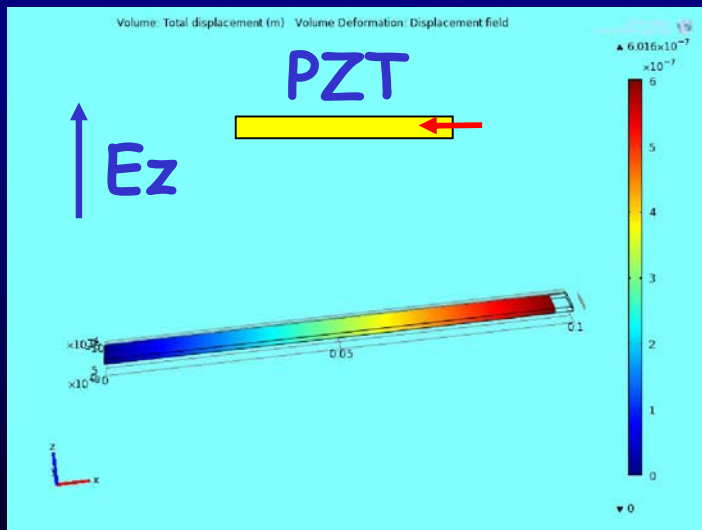
# Principle of PZT inkjet

Use this strain

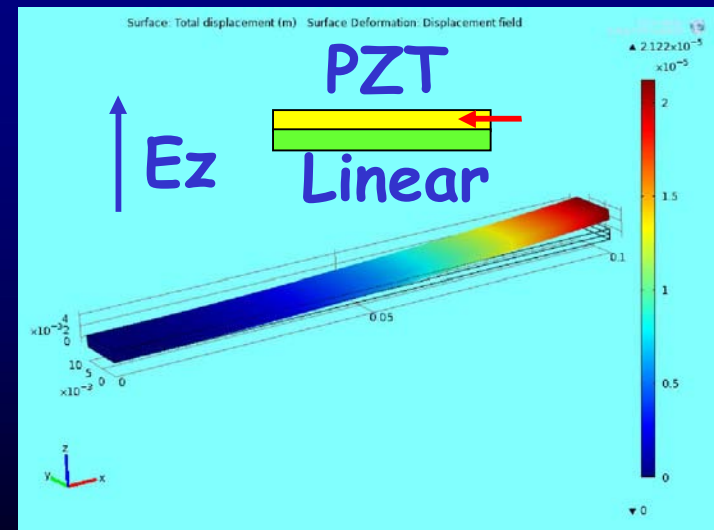
Normal strain

Shear strain

$$\begin{bmatrix} S_{xx} \\ S_{yy} \\ S_{zz} \\ S_{xy} \\ S_{yz} \\ S_{xz} \end{bmatrix} = \begin{bmatrix} 0 & 0 & d_{31} \\ 0 & 0 & d_{32} \\ 0 & 0 & d_{33} \\ 0 & * & 0 \\ * & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} E_x \\ E_y \\ E_z \end{bmatrix}$$



PZT only



Unimorph: PZT+ linear material

# Statement of the problem

## Piezoelectric

$$\frac{\partial^2 U}{\partial t^2} + \nabla \cdot \mathcal{T} \quad \text{in } \Omega_{\text{PZT}}$$

$$\nabla \cdot \mathcal{E} = 0 \quad \text{in } \Omega_{\text{PZT}}$$

$$U = 0 \quad \text{on } \Gamma_{\text{D0PZT}}$$

$$\mathcal{T} \cdot n = 0 \quad \text{on } \Gamma_{\text{N0PZT}}$$

$$\varphi = \varphi_0 \quad \text{on } \Gamma_{\text{D0ELE}}$$

$$\frac{\partial \varphi}{\partial n} = 0 \quad \text{on } \Gamma_{\text{N0ELE}}$$

$$\mathcal{T}_{ij} = c_{ijkl} U_{k,lj} + e_{kij} \varphi_{,k}$$

$$\mathcal{E}_i = -e_{ikl} U_{k,l} + \epsilon_{ik} \varphi_{,k}$$

## Free surface Navier-Stokes

$$\rho \frac{D^2 u}{Dt^2} - \mu \nabla^2 u + \nabla p = f \quad \text{in } \Omega_{\text{ink}}$$

$$\nabla \cdot u = 0 \quad \text{in } \Omega_{\text{ink}}$$

$$u \cdot n = 0 \quad \text{on } \Gamma_{\text{D0ink}}$$

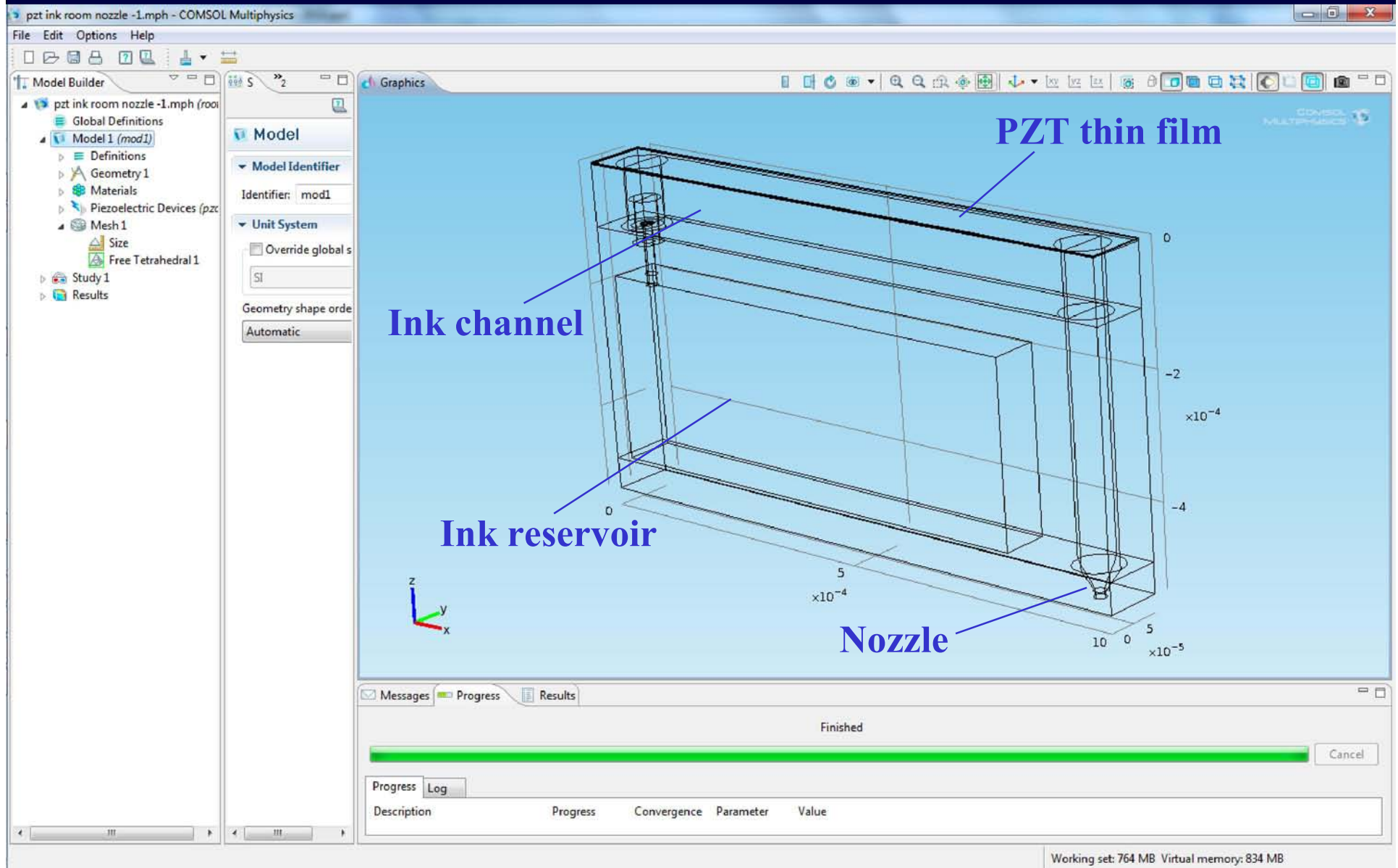
$$p = g \quad \text{on } \Gamma_{\text{Dpink}}$$

$$-\tau \cdot n + \sigma \kappa n = 0 \quad \text{on } \Gamma_{\text{N0ink}}$$

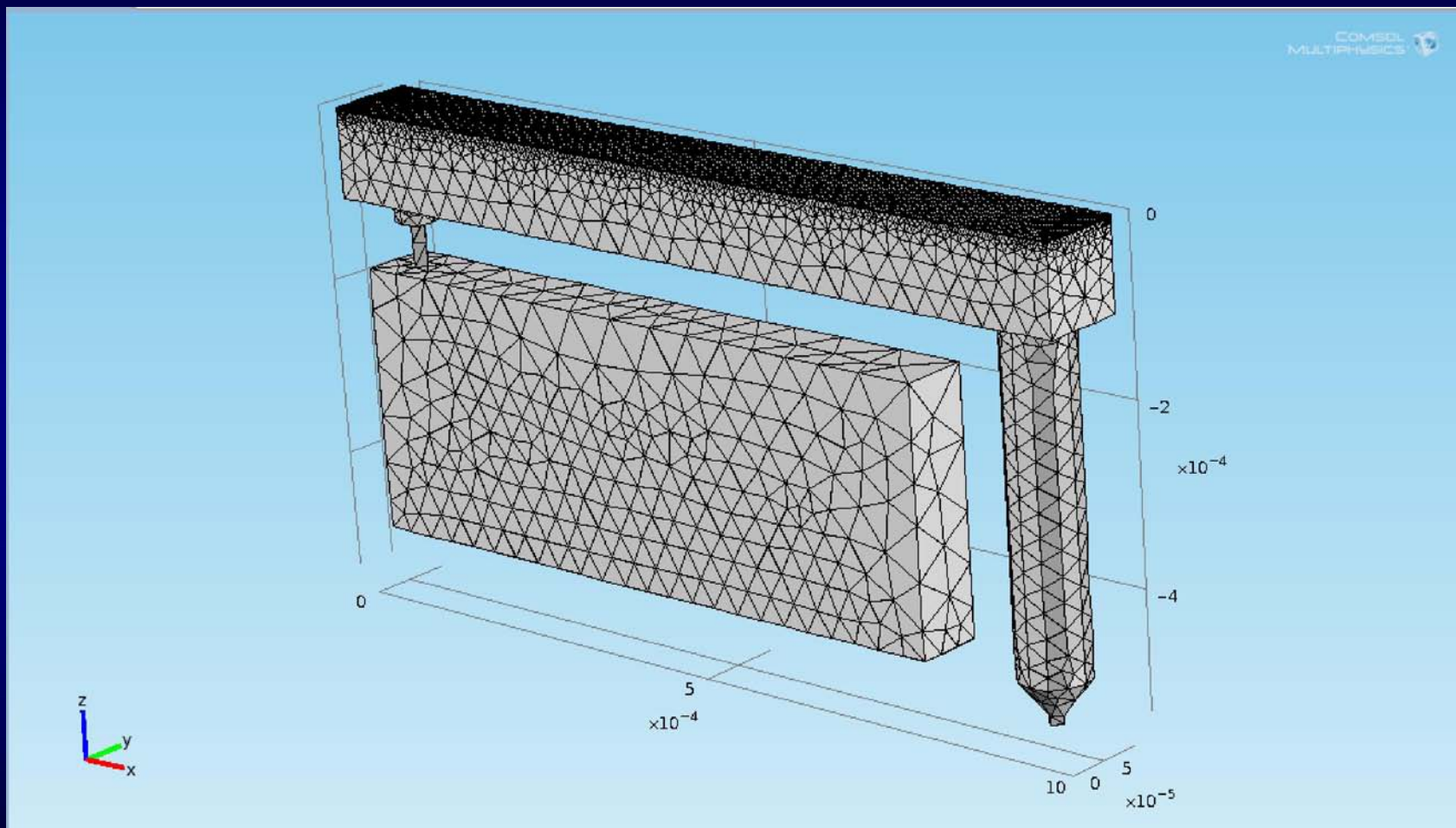
$$\tau_{ij} = \frac{1}{2} (u_{i,j} + u_{j,i})$$

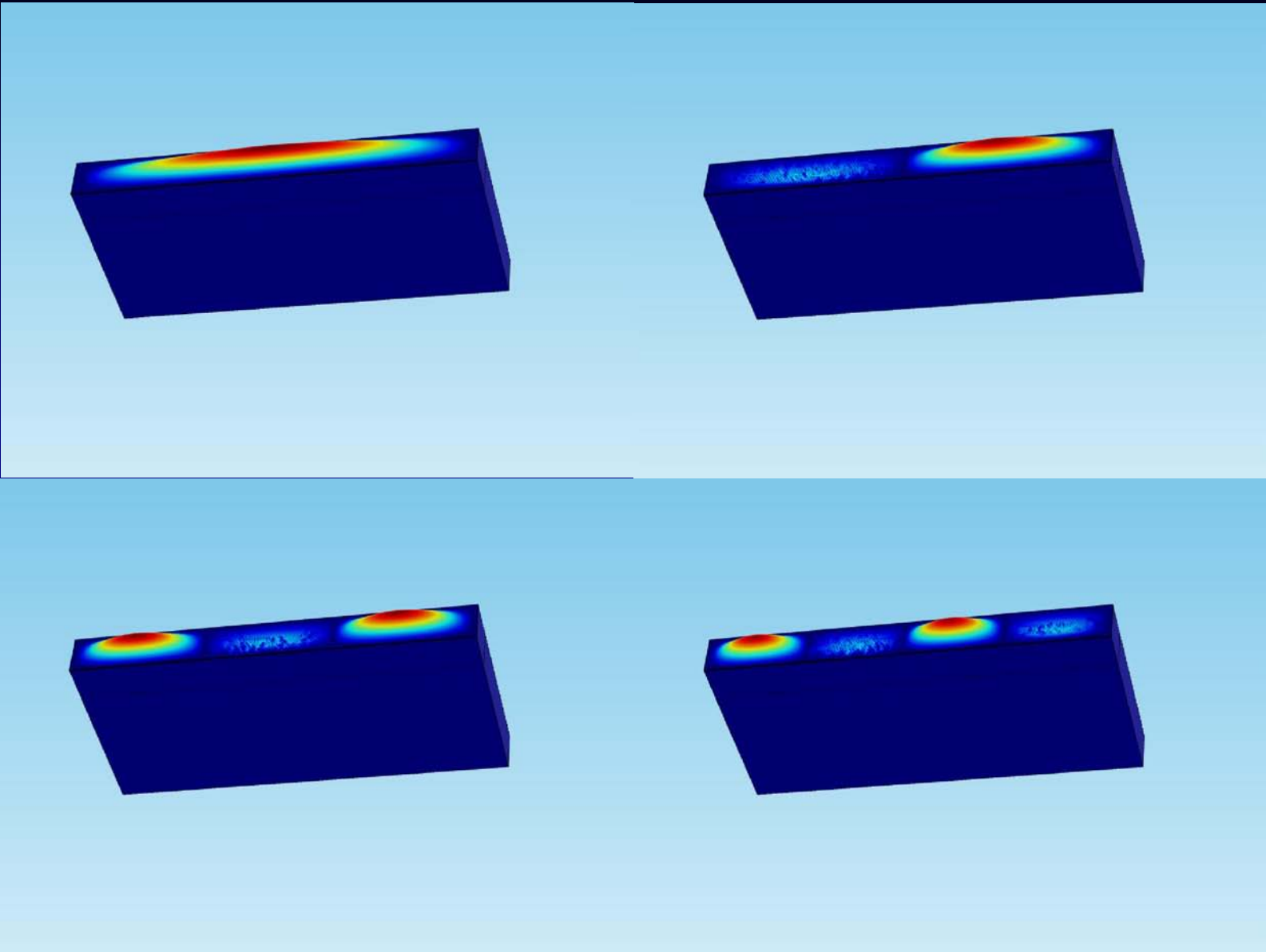
$$\frac{\partial \eta}{\partial t} + u \cdot \nabla \eta = 0 \quad \text{in } \Omega_{\text{ink}}$$

# COMSOL 4.0a modeling



# Industrial inkjet printhead

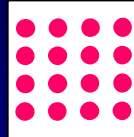
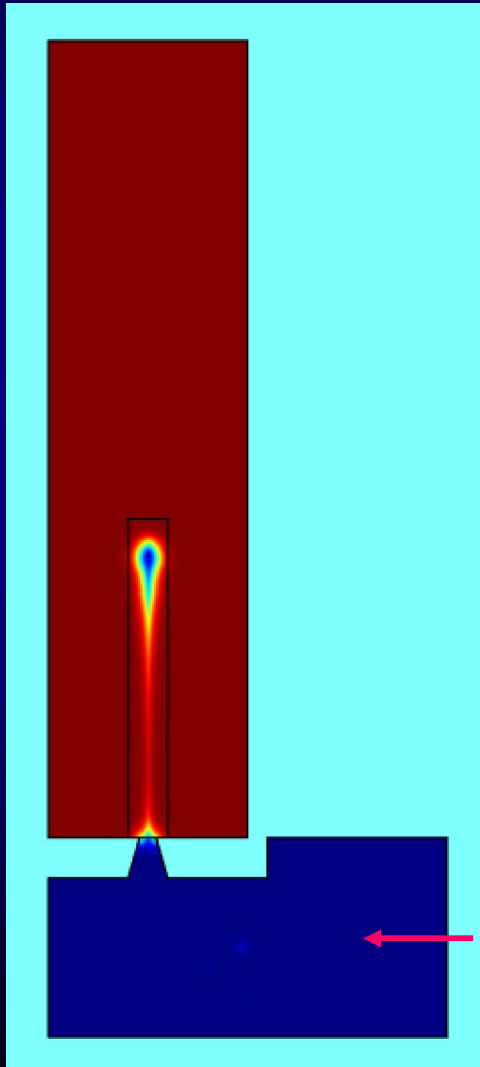




*The first 4 modes around 2MHz*

# Pico-litter droplet formation

## Single pulse

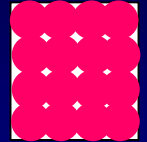
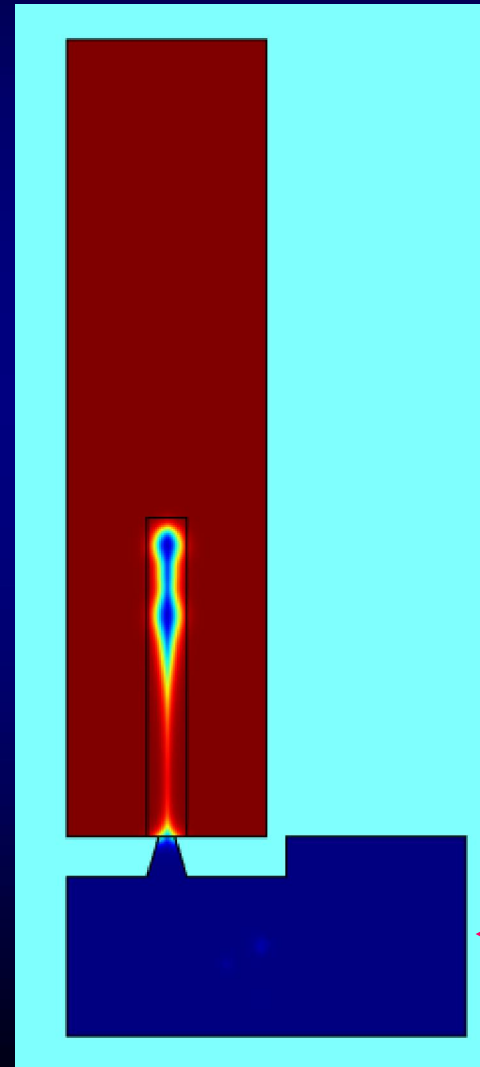


Small dot  
Low density

Pressure



## Double pulse



Large dot  
High density

Pressure



# Surface tension vs pinch off timing

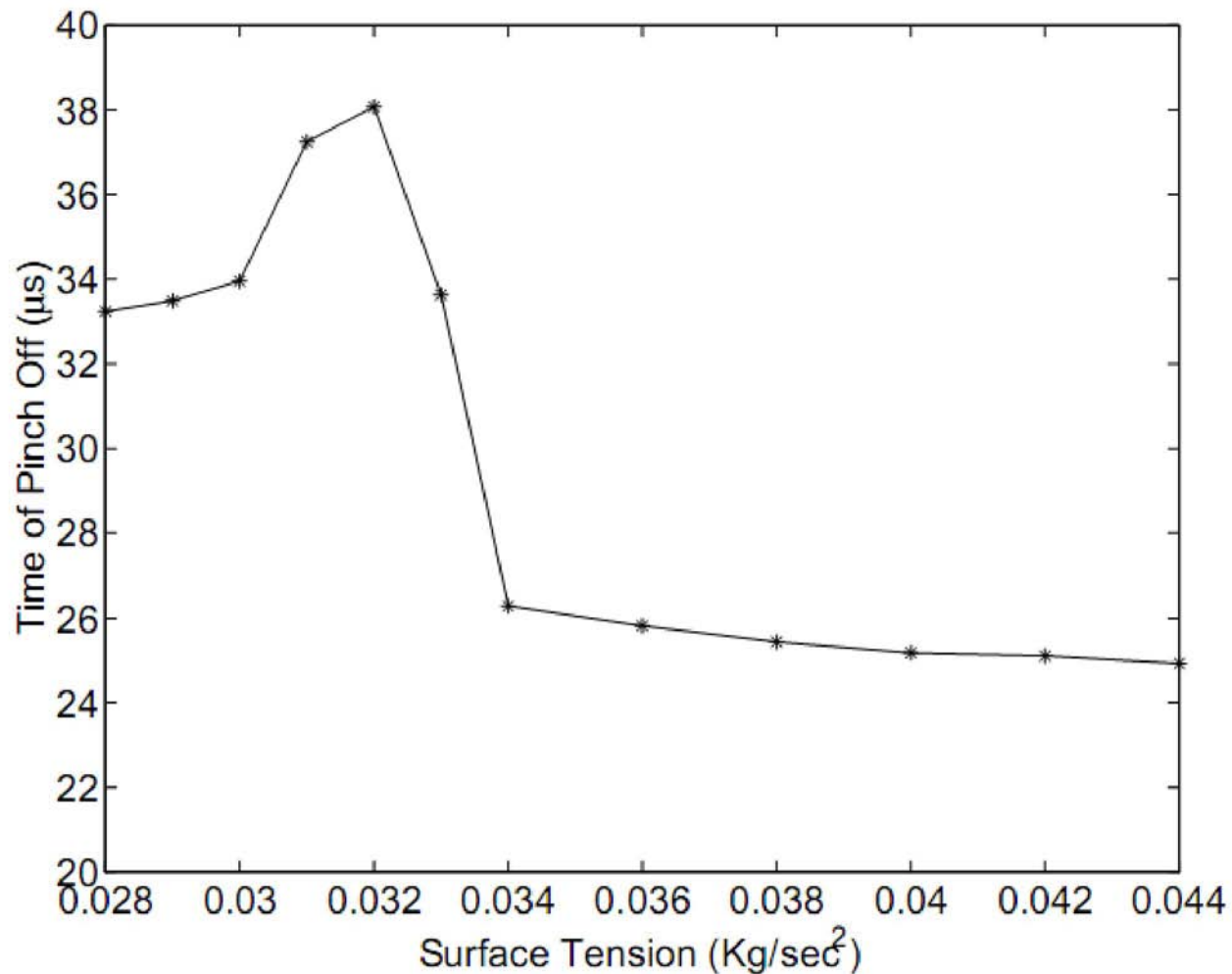


Figure 16: The relation between the time of pinch off and the surface tension.

**J. Yu, S. Sakai, J. A. Sethian, “A coupled levelset projection method applied to inkjet simulation”.**

# Consideration of surface tension

## *Definition of curvature using Levelset function*

$$n = \frac{\nabla\phi}{|\nabla\phi|} \quad \kappa = \nabla \cdot n = \nabla \cdot \left( \frac{\nabla\phi}{|\nabla\phi|} \right) \Rightarrow \text{2nd derivative}$$

coordinates  $(r, z)$ , the curvature can be expanded as

$$\kappa(\phi) = \nabla \cdot \left( \frac{\nabla\phi}{|\nabla\phi|} \right) = \frac{\phi_{rr}\phi_z^2 - 2\phi_{rz}\phi_r\phi_z + \phi_{zz}\phi_r^2 + \phi_r(\phi_r^2 + \phi_z^2)/r}{(\phi_r^2 + \phi_z^2)^{3/2}},$$

**J. Yu, S. Sakai, J. A. Sethian, “A coupled levelset projection method applied to inkjet simulation”.**



# Consideration of surface tension

## Surface tension term in COMSOL (from manual)

### Force Terms

The four forces on the right-hand side of Equation 0-1 are due to gravity, surface tension, a force due to an external contribution to the free energy (using the phase field method only), and a user defined volume force.

- The surface tension force for the level set method acting at the interface between the two fluids is

$$\mathbf{F}_{st} = \sigma \kappa \delta \mathbf{n}$$

where  $\sigma$  is the surface tensions coefficient (SI unit: N/m),  $\kappa$  is the curvature, and  $\mathbf{n}$  is the unit normal to the interface, as defined in [Variables For Geometric Properties of the Interface](#).  $\delta$  (SI unit: 1/m) is a Dirac delta function

concentrated to the interface.  $\kappa$  depends on second derivatives of the level set function  $\phi$ . This can lead to poor accuracy of the surface tension force. Therefore, the program uses the alternative formulation

$$\mathbf{F}_{st} = \nabla \cdot (\sigma(\mathbf{I} - (\mathbf{n}\mathbf{n}^T))\delta)$$

### Conservative body force

In the weak formulation of the momentum (fluid-flow) equations, it is possible to move the divergence operator, using integration by parts, to the test functions for the velocity components.

$$- \int_{\Omega} (\sigma \kappa n \delta) v \, dx = \int_{\Omega} \sigma (I - (n n^T)) \delta \nabla v \, dx$$

Weak form

$$- \int_{\Gamma} \sigma (I - (n n^T)) \delta v \, dx$$

# Consideration of surface tension

## *Surface tension as surface force by Melcher, Bansch*

in (1) ). If  $\varphi$  is a smooth vector valued function on  $\Gamma$  we can integrate by parts to get

### Surface force

$$\int_{\Gamma} \kappa \nu \cdot \varphi = \int_{\Gamma} (\underline{\Delta} id_{\Gamma}) \cdot \varphi = - \int_{\Gamma} \underline{\nabla} id_{\Gamma} \cdot \underline{\nabla} \varphi.$$

and the Laplace–Beltrami operator

$$\underline{\Delta} f := \frac{1}{\sqrt{\det g}} \partial_{u_i} \left( \sqrt{\det g} g^{ij} \partial_{u_j} (f \circ \chi) \right),$$

see for instance [11, 18]. We will make use of the identity

$$\underline{\Delta} id_{\Gamma} = \kappa \nu$$

**Eberhard Bansch, “Finite element discretization of the Navier-Stokes equations with a free capillary surface”, *Numer. Math.* (2001)**

# Consideration of surface tension

*Surface tension as surface force by Mizuyama*

$$\kappa = \frac{\ddot{x}\dot{y} - \dot{x}\ddot{y}}{(\dot{x}^2 + \dot{y}^2)^{3/2}}$$

Definition from  
differential geometry

$$\kappa n_i = -\frac{d^2 x_i}{ds^2}$$

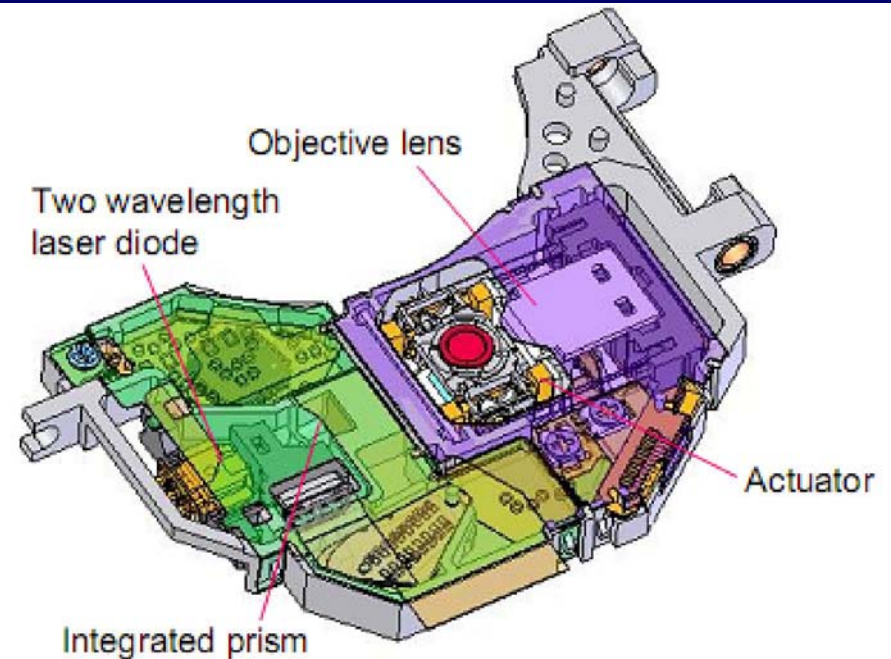
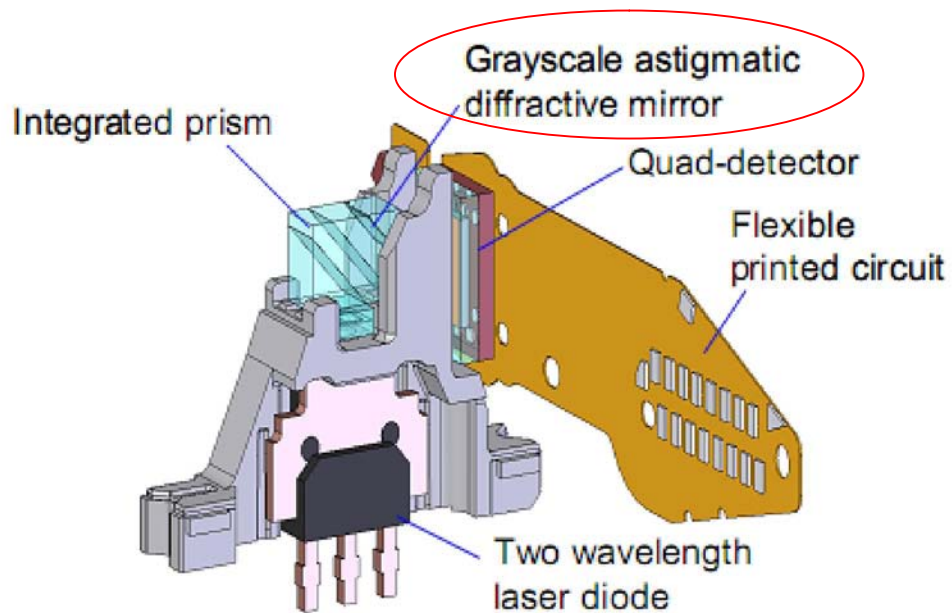
s: Arc length

$$\int_{\Gamma} v \sigma \kappa n ds = -\sigma \int_{\Gamma} v \frac{d^2 x}{ds^2} ds$$

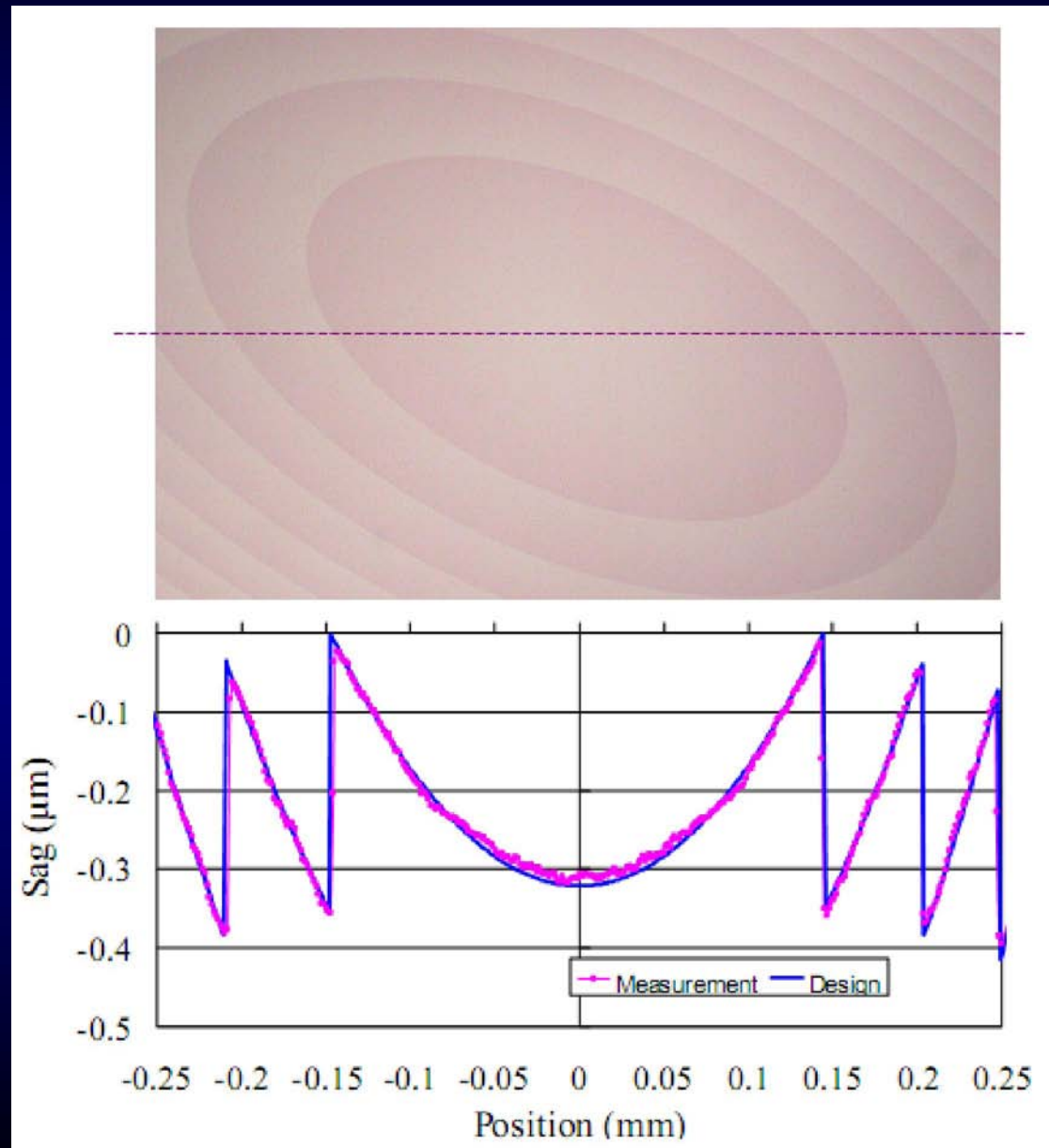
Surface force in  
weak form

$$= \sigma \int_{\Gamma} \frac{dv}{ds} \frac{dx}{ds} ds - \sigma \left[ v \frac{dx}{ds} \right]_{\partial \Gamma}$$

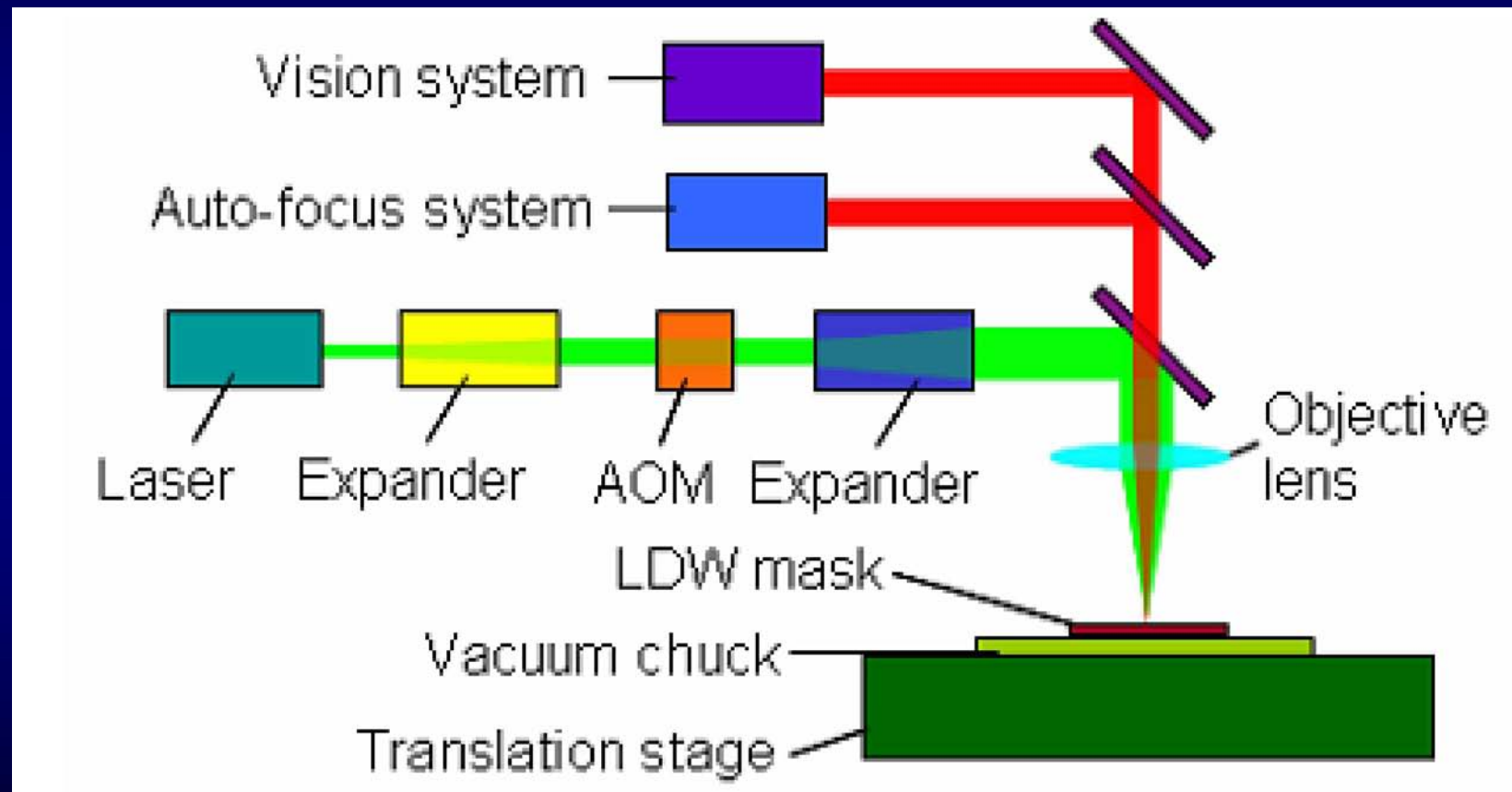
# Grayscale photolithography technology for optical pickup



# Grayscale mask

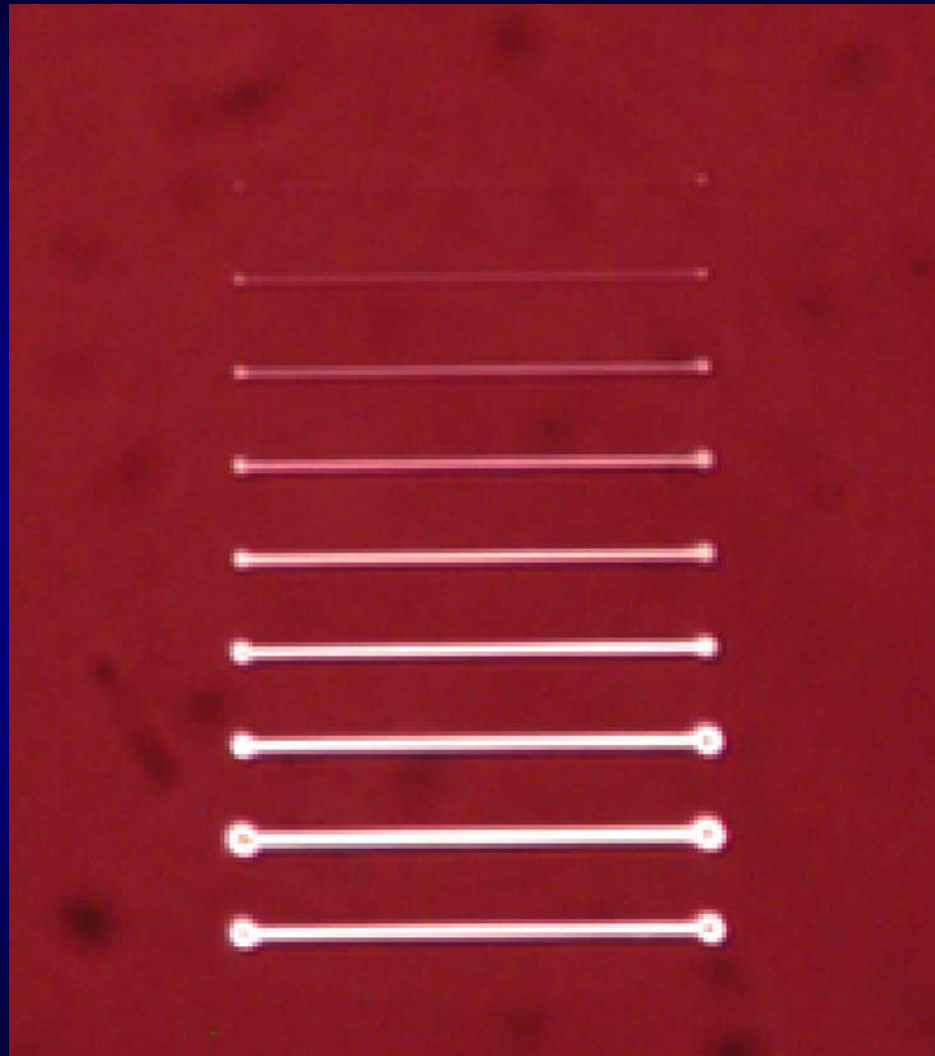


# Laser direct write



# Laser direct write

Laser power



20%

30%

40%

50%

60%

70%

80%

90%

100%

Transmission changes as laser power

# Statement of the problem

$$\left\{ \begin{array}{ll} \rho c_p \frac{\partial T}{\partial t} - \nabla \cdot (\kappa \nabla T) = 0 & \text{in } \Omega \\ \kappa \nabla T \cdot n = I(x, t) & \text{on } \Gamma_N \quad \text{Surface heat source} \\ \kappa \nabla T \cdot n = 0 & \text{on } \Gamma_{N0} \\ T = T_0 & \text{on } \Gamma_D \end{array} \right.$$
$$I(x, t) = I_0(x, t) (1 - \tau(x, t) - r_0) \quad \text{Transmission}$$

## Gaussian beam

$$I_0(x, t) = \exp \left\{ -\frac{(x - x_c - vt)^2}{w_x^2} - \frac{(y - y_c)^2}{w_y^2} \right\}$$

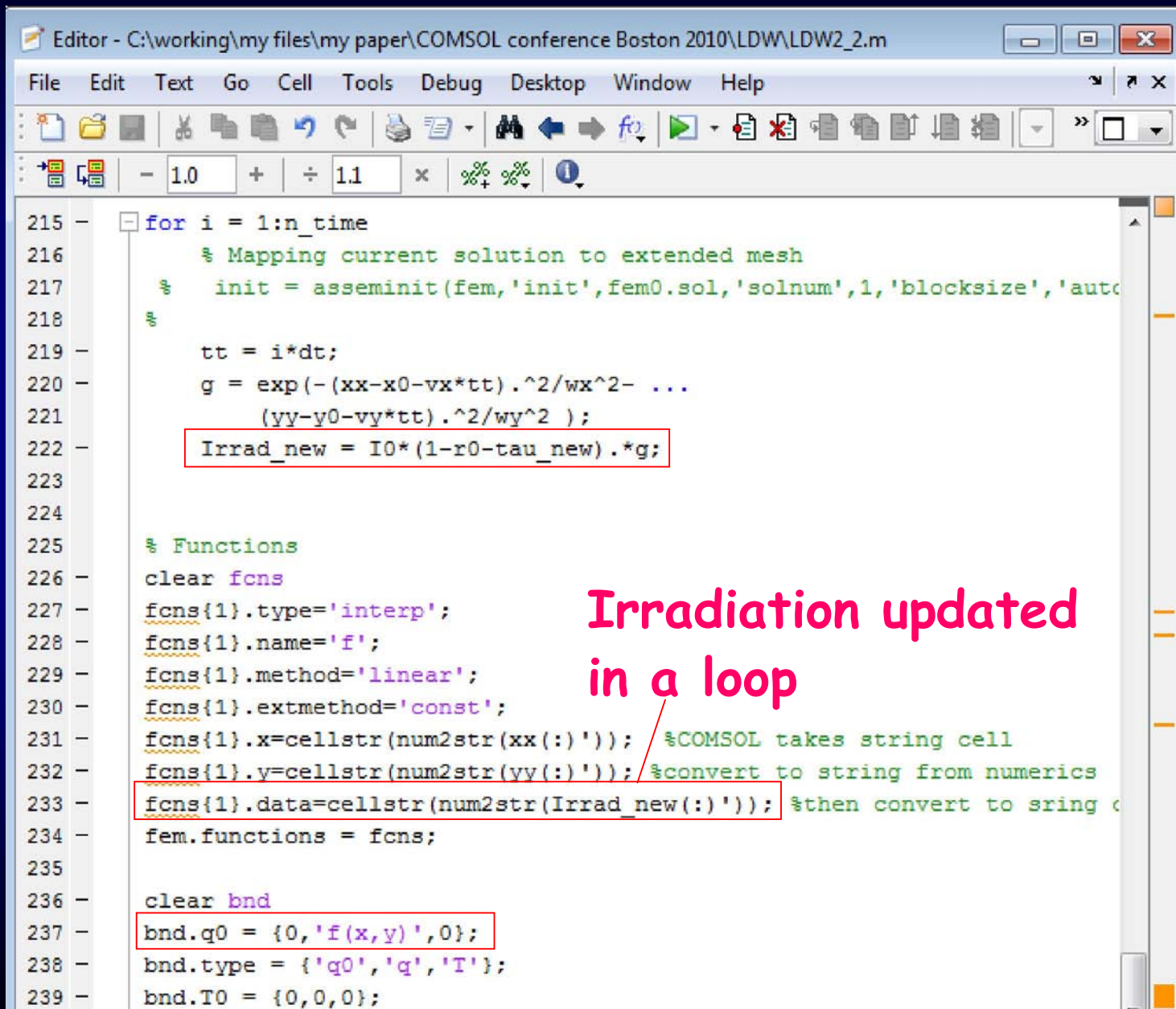
## Transmission vs temperature

$$\tau(\bar{T}) = \tau_{min} + (\tau_{max} - \tau_{min}) \sin^2(\pi \bar{T} / 2T_0) \quad T < T_0$$

$$\bar{T}(x, t) = \max\{T(x, t'), t' < t\} \quad \text{Irreversible process}$$



# COMSOL with MATLAB

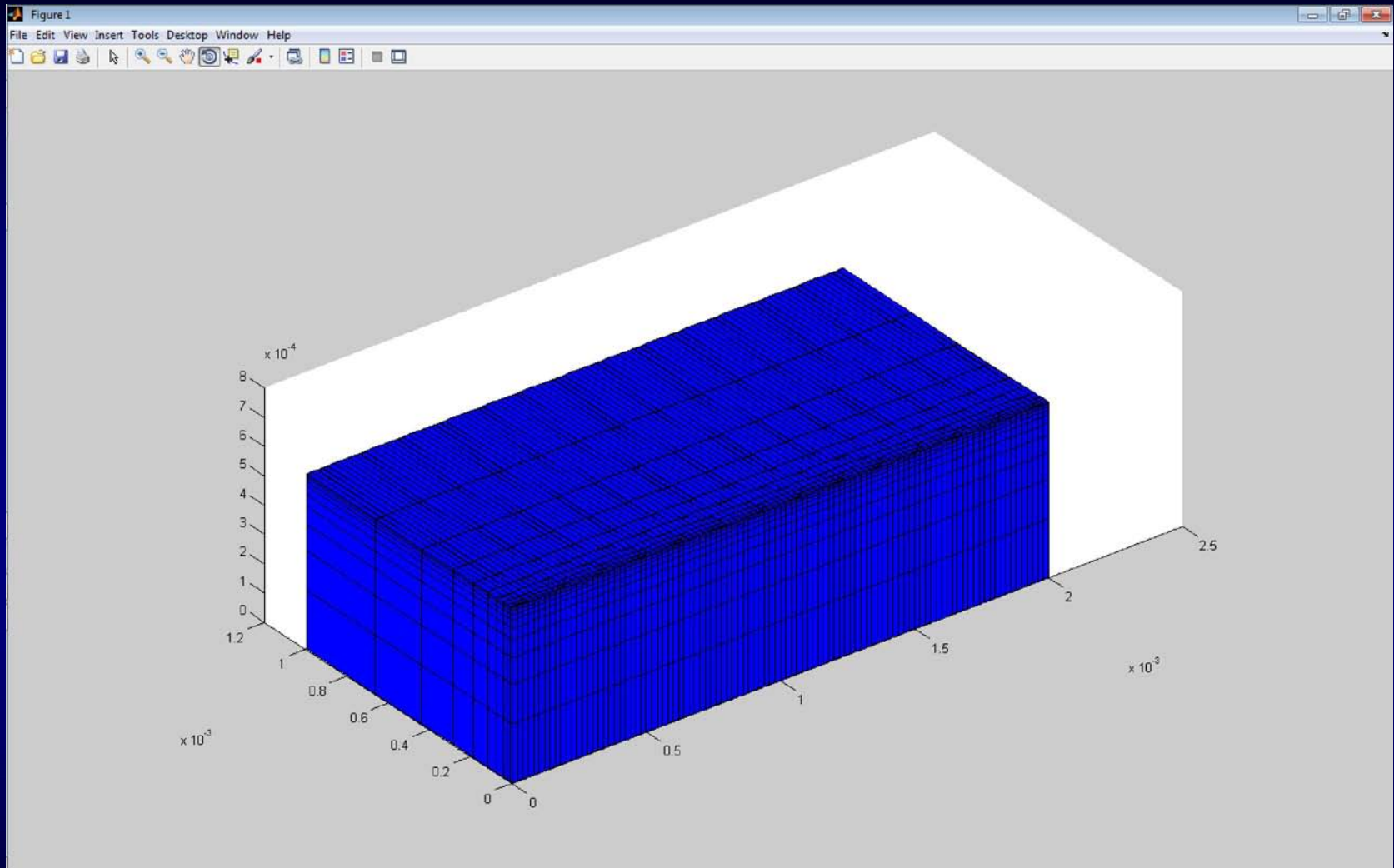


The image shows a MATLAB Editor window titled "Editor - C:\working\my files\my paper\COMSOL conference Boston 2010\LDW\LDW2\_2.m". The code is a MATLAB script for COMSOL, enclosed in a loop. The code is as follows:

```
215 - for i = 1:n_time
216     % Mapping current solution to extended mesh
217     % init = asseminit(fem,'init',fem0.sol,'solnum',1,'blocksize','auto
218     %
219     tt = i*dt;
220     g = exp(-(xx-x0-vx*tt).^2/wx^2- ...
221         (yy-y0-vy*tt).^2/wy^2 );
222     Irrad_new = I0*(1-r0-tau_new).*g;
223
224
225     % Functions
226     clear fcns
227     fcns{1}.type='interp';
228     fcns{1}.name='f';
229     fcns{1}.method='linear';
230     fcns{1}.extmethod='const';
231     fcns{1}.x=cellstr(num2str(xx(:)')); %COMSOL takes string cell
232     fcns{1}.y=cellstr(num2str(yy(:)')); %convert to string from numerics
233     fcns{1}.data=cellstr(num2str(Irrad_new(:)')); %then convert to string c
234     fem.functions = fcns;
235
236     clear bnd
237     bnd.q0 = {0,'f(x,y)',0};
238     bnd.type = {'q0','q','T'};
239     bnd.T0 = {0,0,0};
```

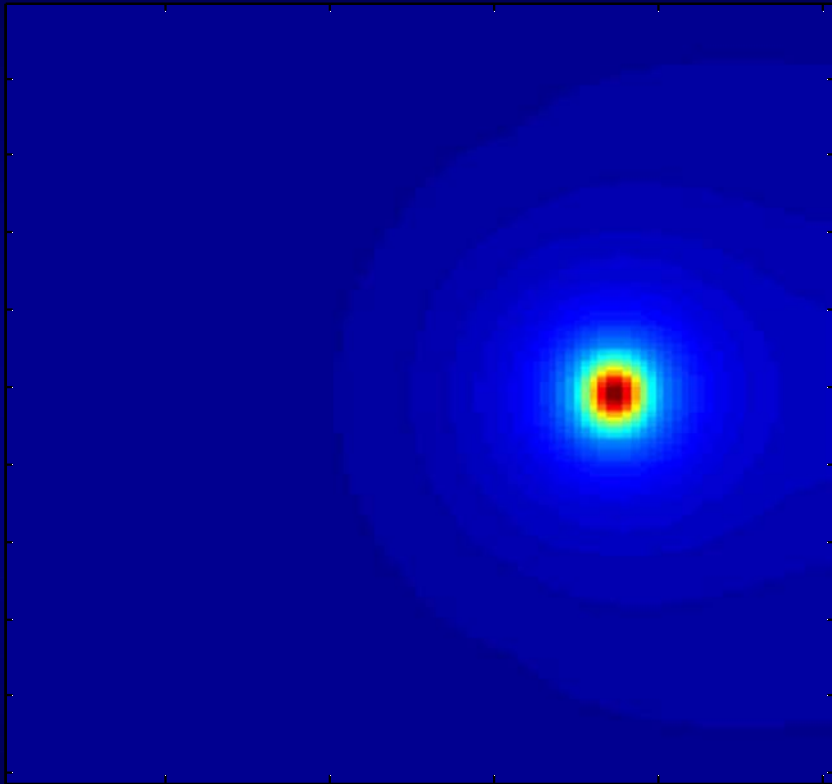
Annotations in the image include:

- A red box around line 222: `Irrad_new = I0*(1-r0-tau_new).*g;`
- A red box around line 233: `fcns{1}.data=cellstr(num2str(Irrad_new(:)')); %then convert to string c`
- A red box around line 237: `bnd.q0 = {0,'f(x,y)',0};`
- A red arrow pointing from the text "Irradiation updated in a loop" to the line 222 box.

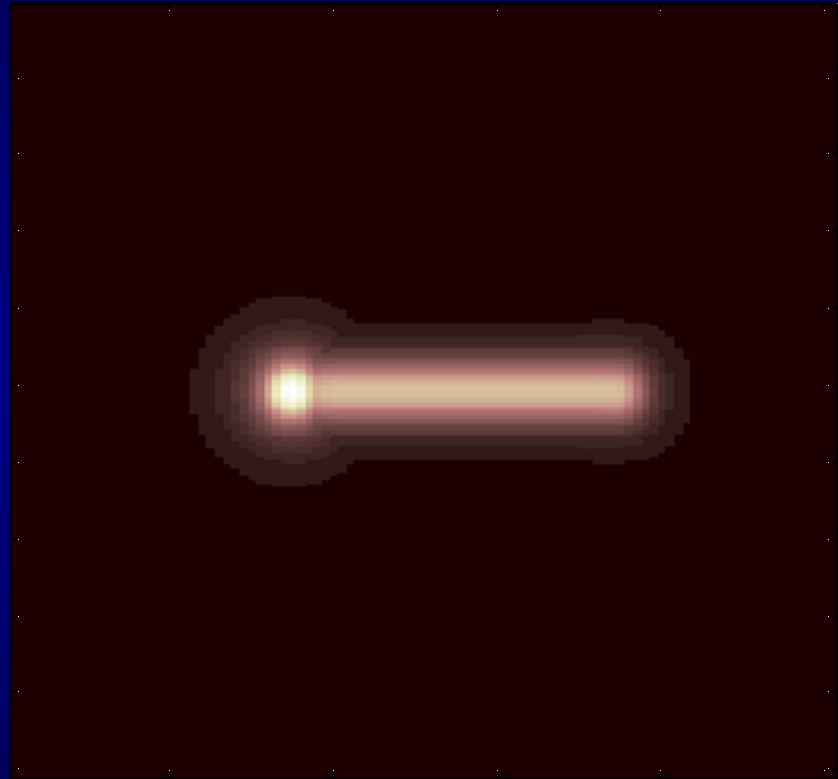


Mesh in 1/4 model

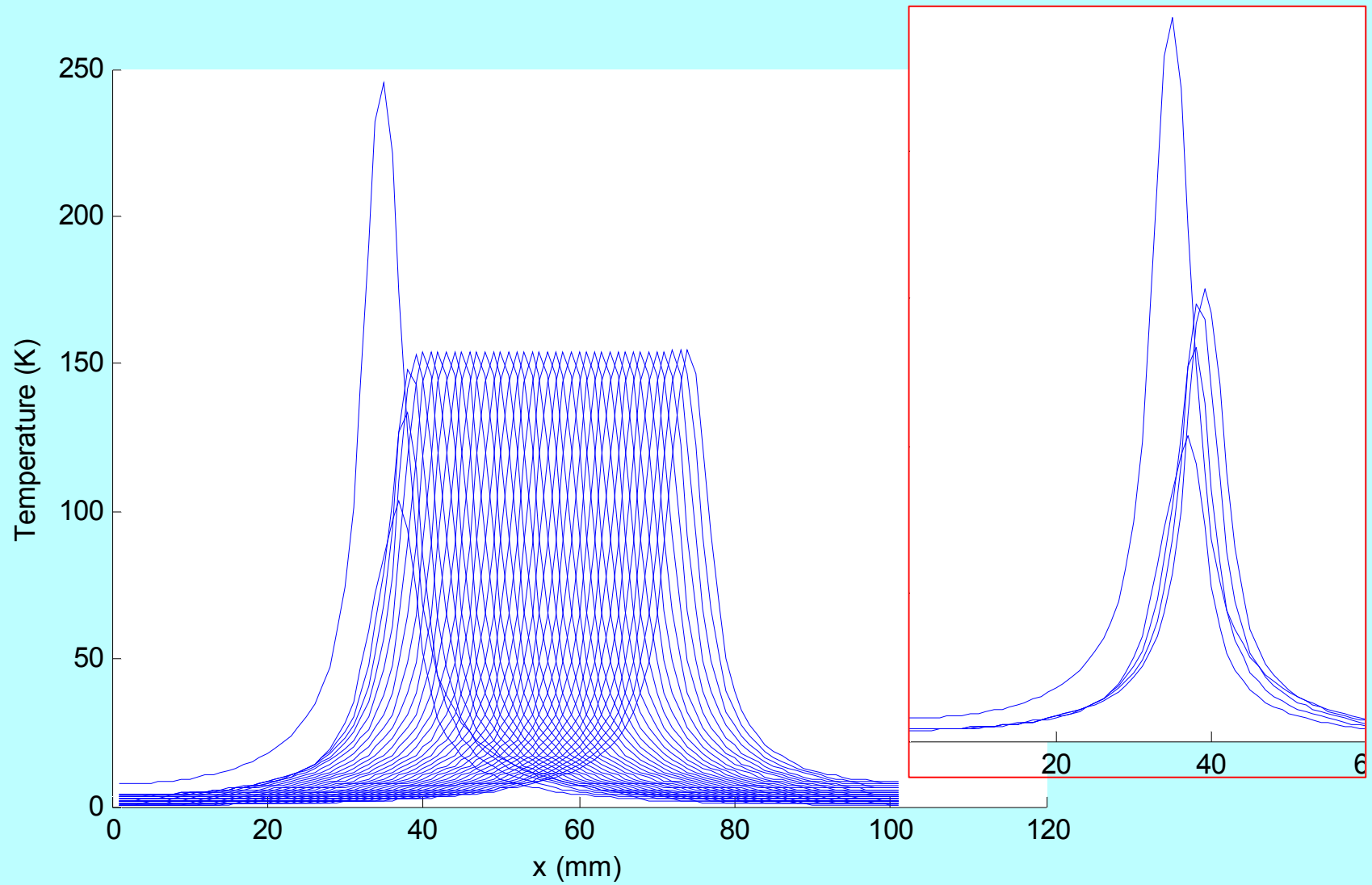
# COMSOL+MATLAB simulation



Temperature



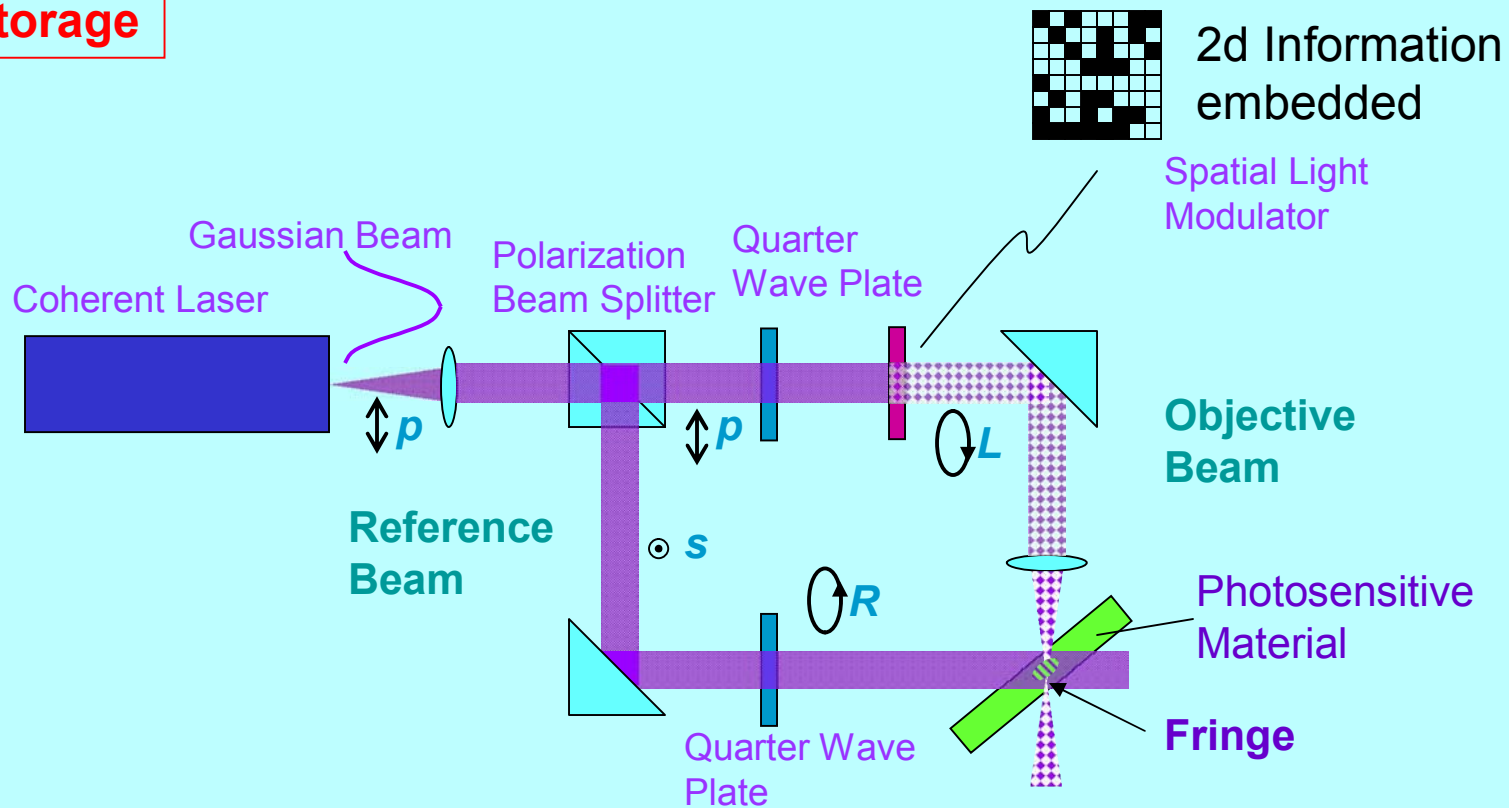
Transmittance



Temperature transition

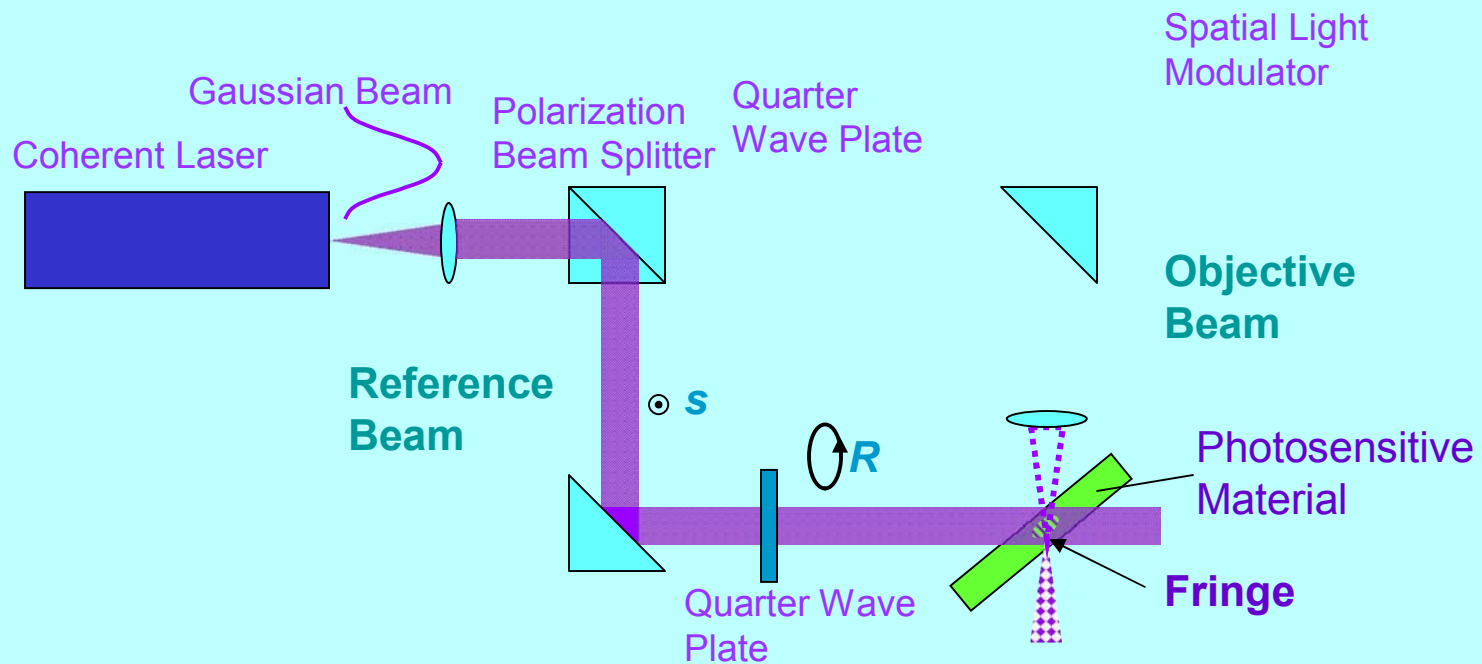
# Holographic data storage

**Storage**

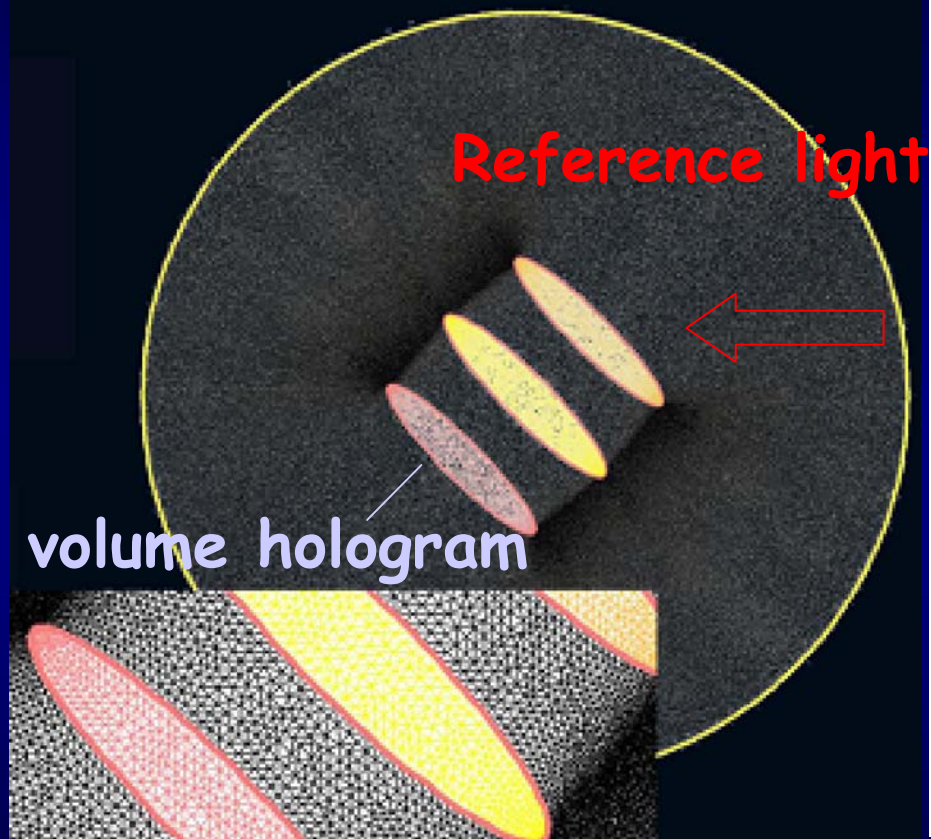


# Holographic data storage

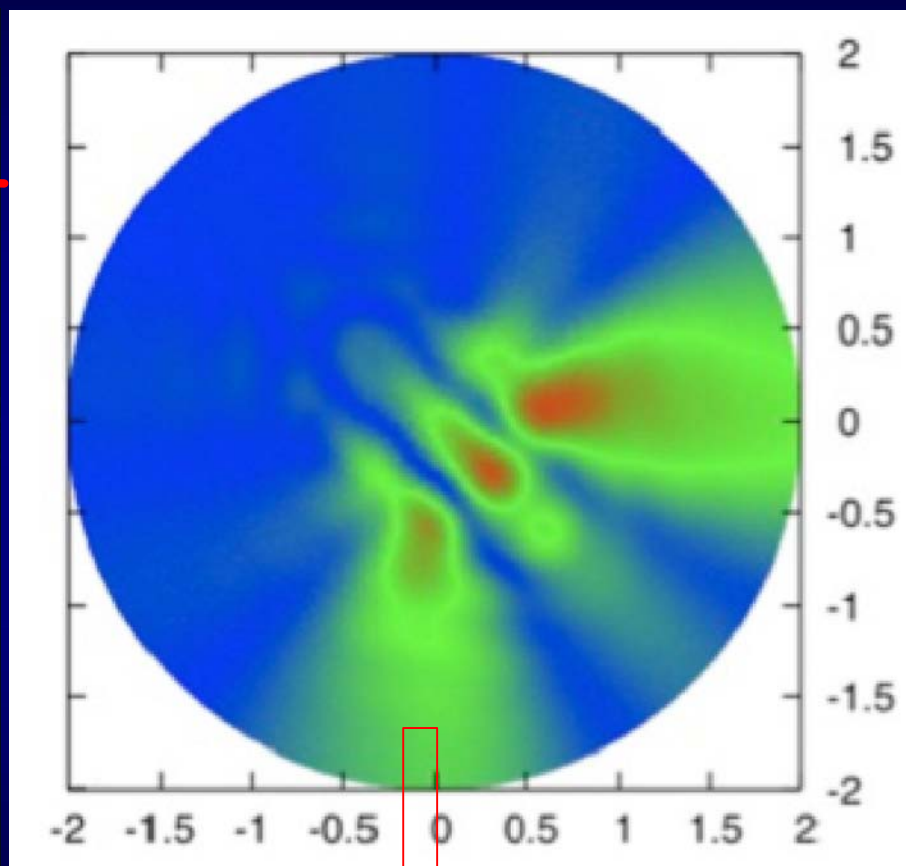
## Retrieval



# Holographic data storage



Fringes in a volume hologram



Information light retrieved

# Statement of the problem

## Infinite region problem

$$-\nabla^2 u - k_2^2 u = (\nabla^2 + k_2^2)u^{inc} \quad \text{in } \Omega_1$$

$$-\nabla^2 u - k_1^2 u = 0 \quad \text{in } \Omega_\infty$$

$$[u]_2^1 = 0 \quad \text{on } \Gamma_{12}$$

$$\left[ \frac{\partial u}{\partial n} \right]_2^1 = 0 \quad \text{on } \Gamma_{12}$$

**Sommerfeld  
radiation condition**

$$\lim_{r \rightarrow +\infty} \sqrt{r} \left( \frac{\partial u}{\partial n} - ik_1 u \right) = 0$$



# Equivalent problem with DtN map

## Finite region problem

$$-\nabla^2 u - k_2^2 u = (\nabla^2 + k_2^2) u^{inc} \quad \text{in } \Omega_1$$

$$-\nabla^2 u - k_1^2 u = 0 \quad \text{in } \Omega_2$$

$$[u]_2^1 = 0 \quad \text{on } \Gamma_{12}$$

$$\left[ \frac{\partial u}{\partial n} \right]_2^1 = 0 \quad \text{on } \Gamma_{12}$$

$$\frac{\partial u}{\partial n} = -\mathcal{S}u \quad \text{on } \Gamma_2$$

**Steklov-Poincare  
operator**

$$\mathcal{S}u = -k_1 \sum_{n=-\infty}^{+\infty} \frac{H_n^{(1)'}(k_1 a)}{H_n^{(1)}(k_1 a)} u_n(a) \phi_n(\theta)$$

# Weak form

$$a(u, v) + s(u, v) = \langle f, v \rangle \quad \forall u \in V$$

$$a(u, v) = \int_{\Omega} (\nabla u \cdot \nabla \bar{v} - k^2 u \bar{v}) dx$$

$$s(u, v) = \int_{\Gamma_2} (\mathcal{S}u) \bar{v} ds$$

$$\langle f, v \rangle := \int_{\Omega} f v dx$$

$$f(x) := \begin{cases} 0 & \text{in } \Omega_2 \\ (\Delta + k_2^2)u^{\text{inc}} & \text{in } \Omega_1 \end{cases}$$

$$s(u, v) = -k_1 a \sum_{n=-\infty}^{+\infty} \frac{H_n^{(1)'}(k_1 a)}{H_n^{(1)}(k_1 a)} u_n(a) \bar{v}_n$$

Thank you!